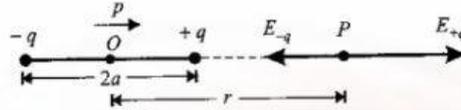


## Unit-1 Electrostatics

### (Electric Charges and Fields)

#### 1. Derive the expression for electric field intensity at a point on the axial line of a dipole.

As shown in the figure, consider an electric dipole of charges  $-q$  and  $+q$ , separated by distance  $2a$  and placed in vacuum. Let  $P$  be a point on the axial line at a distance  $r$  from the centre  $O$  of the dipole on the side of the charge  $+q$ .



Electric field due to charge  $-q$  at a point  $P$  is

$$\vec{E}_{-q} = \frac{-q}{4\pi\epsilon_0(\quad)^2} \hat{p}$$

Where  $\hat{p}$  is a unit vector along the dipole axis from  $-q$  to  $+q$ .

Electric field due to charge  $+q$  at a point  $P$  is

$$\vec{E}_{+q} = \frac{q}{4\pi\epsilon_0(\quad)^2} \hat{p}$$

Hence the resultant electric field at point  $P$  is

$$\begin{aligned} \vec{E}_{axial} &= \vec{E}_{-q} + \vec{E}_{+q} \\ &= \frac{q}{4\pi\epsilon_0} \left[ \frac{1}{(\quad)^2} - \frac{1}{(\quad)^2} \right] \hat{p} = \frac{q}{4\pi\epsilon_0} \cdot \frac{\quad}{(\quad^2 - \quad^2)^2} \hat{p} \end{aligned}$$

$$\boxed{\vec{E}_{axial} = \frac{1}{4\pi\epsilon_0} \cdot \frac{\quad}{(\quad^2 - \quad^2)^2} \hat{p}}$$

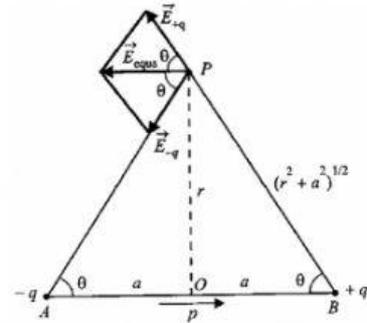
Here  $p = q \times 2a =$  dipole moment. For  $r \gg a$ ,  $a^2$  can be neglected compared to  $r^2$ .

$$\therefore \boxed{\vec{E}_{axial} = \frac{1}{4\pi\epsilon_0} \cdot \frac{\quad}{\quad^3} \hat{p}}$$

Electric field at any point on the axis of the dipole acts along the dipole axis from \_\_\_\_\_ to \_\_\_\_\_ charge, i.e. in the direction of \_\_\_\_\_.

**2. Derive the expression for the electric field intensity at any point along the equatorial line of an electric dipole.**

As shown in the figure, consider an electric dipole of charges  $-q$  and  $+q$ , separated by distance  $2a$  and placed in vacuum. Let  $P$  be point on the equatorial line of the dipole at distance  $r$  from it. i.e.  $OP = r$



Electric field at point  $P$  due to  $+q$  charge is

$$\vec{E}_{+q} = \frac{q}{4\pi\epsilon_0(\underline{\quad}^2 + \underline{\quad}^2)} \text{ directed along } BP.$$

Electric field at point  $P$  due to  $-q$  charge is

$$\vec{E}_{-q} = \frac{q}{4\pi\epsilon_0(\underline{\quad}^2 + \underline{\quad}^2)} \text{ directed along } AP.$$

Clearly, the magnitudes of  $\vec{E}_{+q}$  and  $\vec{E}_{-q}$  are equal. So, the components of  $\vec{E}_{+q}$  and  $\vec{E}_{-q}$  normal to the dipole axis will cancel out. The components parallel to the dipole axis will add up. The total electric field  $\vec{E}_{eq}$  is opposite to  $\vec{p}$ .

$$\begin{aligned} \therefore \vec{E}_{eq} &= -(E_{-q} \underline{\quad} \theta + E_{+q} \underline{\quad} \theta) \hat{p} \\ \vec{E}_{eq} &= -\underline{\quad} \vec{E}_{+q} \underline{\quad} \theta \hat{p} \quad (\because E_{+q} = E_{-q}) \\ \text{or, } \vec{E}_{eq} &= -2 \frac{q}{4\pi\epsilon_0(\underline{\quad}^2 + \underline{\quad}^2)} \underline{\quad} \theta \hat{p} \end{aligned}$$

Now, from fig.  $\theta = \frac{a}{\sqrt{r^2 + a^2}}$ ,

$$\therefore \vec{E}_{eq} = -2 \frac{q}{4\pi\epsilon_0(\underline{\quad}^2 + \underline{\quad}^2)} \frac{a}{\sqrt{r^2 + a^2}} \hat{p}$$

$$\text{or, } \boxed{\vec{E}_{eq} = -\frac{1}{4\pi\epsilon_0} \cdot \frac{\underline{\quad}}{(\underline{\quad}^2 + \underline{\quad}^2)^{3/2}} \hat{p}} \text{ where } \underline{\quad} = 2qa, \text{ is electric dipole moment}$$

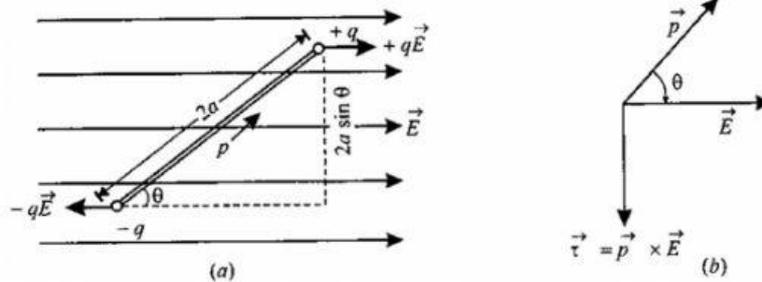
If the point  $P$  is located far away from the dipole,  $r \gg a$ , then

$$\boxed{\vec{E}_{eq} = -\frac{1}{4\pi\epsilon_0} \cdot \frac{\underline{\quad}}{\underline{\quad}^3} \hat{p}}$$

The direction of electric field at any point on the equatorial line of the dipole will be antiparallel to the dipole moment  $\hat{p}$ .

The electric field intensity due to a short dipole at a distance ' $r$ ' along its axis is  $\underline{\quad}$  the intensity at the same distance along the equatorial axis.

3. Derive the expression for the torque acting on an electric dipole, when held in a uniform electric field. (net translational force is zero)



Consider an electric dipole consisting of charges  $+q$  and  $-q$  and length  $2a$  placed in a uniform electric field  $\vec{E}$ , making an angle  $\theta$  with it. It has a dipole moment of magnitude,

$$p = \underline{\hspace{2cm}}$$

Force exerted on charge  $+q$  by field  $\vec{E} = +\underline{\hspace{1cm}}\vec{E}$  along  $\vec{E}$

Force exerted on charge  $-q$  by field  $\vec{E} = -\underline{\hspace{1cm}}\vec{E}$  opposite to  $\vec{E}$

$$\therefore \vec{F}_{total} = +\underline{\hspace{1cm}}\vec{E} - \underline{\hspace{1cm}}\vec{E} = \underline{\hspace{2cm}}$$

Hence net translating force on dipole in a uniform electric field is  $\underline{\hspace{2cm}}$ . But the two equal and opposite force act at different points on the dipole. So, they form a couple and exert  $\underline{\hspace{2cm}}$ .

$\underline{\hspace{2cm}}$  = either force x perpendicular distance between the two forces

$$i.e. \tau = \underline{\hspace{1cm}}E \times 2a \underline{\hspace{1cm}} \theta = (\underline{\hspace{1cm}} \times 2a)E \underline{\hspace{1cm}} \theta$$

$$\boxed{\tau = \underline{\hspace{1cm}}E \underline{\hspace{1cm}} \theta}$$

As the direction of torque  $\vec{\tau}$  is perpendicular to both  $\vec{p}$  and  $\vec{E}$ , we can write

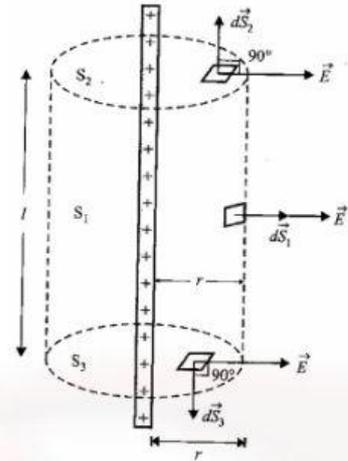
$$\boxed{\vec{\tau} = \underline{\hspace{1cm}} \times \underline{\hspace{1cm}}}$$

As shown in the figure the direction of torque is that in which a right-handed screw would advance when rotated from  $\underline{\hspace{1cm}}$  to  $\underline{\hspace{1cm}}$ .

4. Derive the expression for the electric field intensity due to infinitely long, straight wire of linear charge density  $\lambda \text{ Cm}^{-1}$ .

Consider a thin infinitely long straight wire having a uniform \_\_\_\_\_ charge density  $\lambda \text{ Cm}^{-1}$ . By symmetry, the field  $\vec{E}$  of the line charge is directed radially \_\_\_\_\_ and its magnitude is same at all points equidistant from the line charge. Here we choose a \_\_\_\_\_ Gaussian surface of radius  $r$ , length  $l$  and with its axis along the line charge.

As shown in the figure the Gaussian surface has curved surfaces  $S_1$  and flat circular ends  $S_2$  and  $S_3$ .



Since  $\vec{dS}_1 \parallel \vec{E}$ ,  $\vec{dS}_2 \perp \vec{E}$  and  $\vec{dS}_3 \perp \vec{E}$  only the curved surface contributes towards the total flux.

$$\Phi_E = \oint_S \vec{E} \cdot \vec{ds} = \int_{S_1} \vec{E} \cdot \vec{ds}_1 + \int_{S_2} \vec{E} \cdot \vec{ds}_2 + \int_{S_3} \vec{E} \cdot \vec{ds}_3$$

$$\Phi_E = \int_{S_1} E dS_1 \cos 0^\circ + \int_{S_2} E dS_2 \cos 90^\circ + \int_{S_3} E dS_3 \cos 90^\circ$$

$$\Phi_E = E \int_{S_1} dS_1 = E \times \text{area of curved surface}$$

$$\Phi_E = E \times 2\pi r l$$

Charge enclosed by the Gaussian surface is  $q = \lambda l$

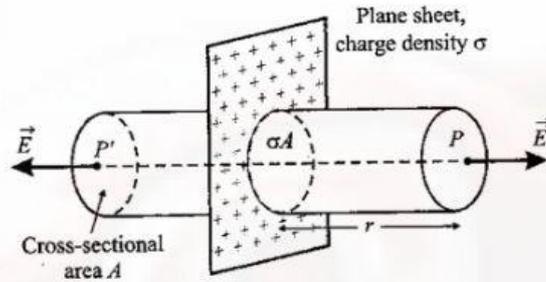
Using Gauss's theorem,  $\Phi_E = \frac{q}{\epsilon_0}$ , we get

$$E \cdot 2\pi r l = \frac{\lambda l}{\epsilon_0}$$

$$\text{or, } \boxed{E = \frac{\lambda}{2\pi\epsilon_0 r}}$$

5. Derive the expression for the electric field intensity at a point near a thin infinite plane sheet of charge density  $\sigma \text{ Cm}^{-2}$ .

As shown in the figure consider a thin, infinite plane sheet of charge with uniform \_\_\_\_\_ charge density  $\sigma$ . We have to calculate its electric field at a point  $P$  at distance  $r$  from it. By symmetry, electric field  $\vec{E}$  points \_\_\_\_\_ normal to the sheet. Also, it must have same magnitude and opposite direction at two points  $P$  and  $P'$  equidistant from the sheet and on opposite sides.



Here we consider a \_\_\_\_\_ Gaussian surface of cross-sectional area  $A$  and length  $2r$  with its axis perpendicular to the sheet. As the lines of force are parallel to the curved surface of the cylinder, the flux through the curved surface is \_\_\_\_\_. The flux through the plane end faces of the cylinder is

$$\phi_E = EA + EA = 2EA$$

Using Gauss's theorem,  $\phi_E = \frac{q}{\epsilon_0}$ , we get

$$2EA = \frac{\sigma A}{\epsilon_0}$$

$$\text{or, } E = \frac{\sigma}{2\epsilon_0}$$

6. Calculate the electric field due to a uniformly charged spherical shell at a point (i) outside the shell, (ii) on the shell and (iii) inside the shell.

Consider a thin spherical shell of charge of radius  $R$  with uniform \_\_\_\_\_ charge density  $\sigma$ . From symmetry, we see that the electric field  $\vec{E}$  at any point is \_\_\_\_\_ and has same magnitude at points equidistant from the centre. To determine electric field at any point at a distance  $r$  from  $O$ , we choose a concentric sphere of radius  $r$  as the \_\_\_\_\_ surface.

(i) When point  $P$  lies outside the spherical shell: The total charge  $q$  inside the Gaussian surface is the charge on the shell of radius  $R$  and area  $4\pi R^2$ .

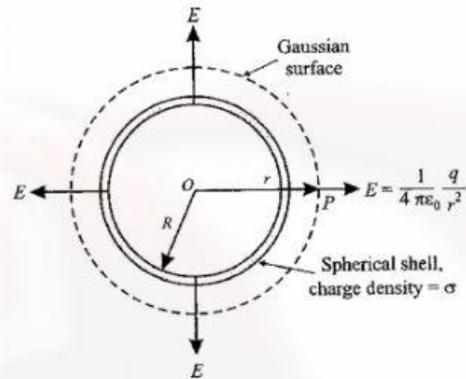
$$\therefore q = 4\pi R^2 \sigma$$

Flux through the Gaussian surface,  $\Phi_E = E \times 4\pi r^2$

Using Gauss's theorem,  $\Phi_E = \frac{q}{\epsilon_0}$ , we get

$$E \times 4\pi r^2 = \frac{q}{\epsilon_0}$$

$$\text{or, } E = \frac{1}{4\pi\epsilon_0} \cdot \frac{q}{r^2}$$



For points outside the shell, the field due to uniformly charged shell is as if the entire charge of shell is concentrated at its \_\_\_\_\_.

(ii) When point  $P$  lies on the spherical shell: The Gaussian surface just encloses the charged spherical shell. Applying Gauss's theorem, for  $r=R$

$$E \times 4\pi R^2 = \frac{q}{\epsilon_0}$$

$$\text{or, } E = \frac{q}{4\pi\epsilon_0 R^2}$$

$$\text{or, } \boxed{E = \frac{\sigma}{\epsilon_0}} \quad \text{since } q = 4\pi R^2 \sigma$$

(iii) When point  $P$  lies inside the spherical shell: For  $r < R$ , the charge enclosed by Gaussian surface is \_\_\_\_\_

Flux through the Gaussian surface,  $\Phi_E = E \times 4\pi r^2$

Using Gauss's theorem,  $\Phi_E = \frac{q}{\epsilon_0}$ , we get

$$E \times 4\pi r^2 = \frac{q}{\epsilon_0} \quad \text{or } E = \frac{q}{4\pi\epsilon_0 r^2}$$

Hence, electric field due to a uniformly charged spherical shell is \_\_\_\_\_ at all points inside the shell.